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Concerted Double Hydrogen-Bond Breaking by Intermolecular Coulombic Decay in the Formic Acid Dimer

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ABSTRACT: Hydrogen bonds are ubiquitous in nature and of fundamental importance to the chemical and physical properties of molecular systems in the condensed phase. Nevertheless, our understanding of the structural and dynamical properties of hydrogenbonded complexes in particular in electronic excited states remains very incomplete. Here, by using formic acid (FA) dimer as a prototype of DNA base pair, we investigate the ultrafast decay process initiated by removal of an electron from the inner-valence shell of the molecule upon electron-beam irradiation. Through fragment-ion and electron coincident momentum measurements and *ab initio* calculations, we find that de-excitation of an outer-valence electron at the same site can initiate ultrafast energy transfer to the neighboring molecule, which is in turn ionized through the emission of low-energy electrons. Our study reveals a concerted breaking of double hydrogen-bond in the dimer initiated by the ultrafast molecular rotations of two FA⁺ cations following this nonlocal decay mechanism.



ntermolecular hydrogen bonds are essential in many areas L of chemistry and biology and govern the properties and dynamics of molecular complexes at the microscopic level.^{1–5} For example, they are the major intermolecular interactions in the liquid and solid structure of water and determine the stability and function of protein and DNA. In the doublestranded DNA, the adenine-thymine and guanine-cytosine base pairs are held together through double and triple hydrogen-bonds, respectively; if these hydrogen bonds are disrupted, it may cause DNA damage. As a result, studying the hydrogen bonds in the dimers or small clusters of molecules has become a subject of great interest for understanding the fundamental features of these intermolecular bonds.⁶⁻⁹ In particular, there has been intense research on the hydrogenbonding dimer of organic formic acid (FA, HCOOH) that is the smallest carboxylic acid and occurs naturally. In the formic acid dimer (FAD), the carboxylic oxygen of each unit is weakly bound to the acidic proton of the other unit via a hydrogenbond, forming an eight-membered ring structure (Figure 1a), which can be regarded as a prototype system for studying the structures and dynamics of hydrogen-bonded base pairs.

In recent years, investigations on the property of excited states in weakly bound systems are attracting considerable interest due to the discovery of a wealth of intermolecular relaxation processes. It was shown that the environment of a molecule with an inner-valence vacancy can actively take part in its de-excitation, giving rise to an ultrafast relaxation mechanism, predicted by Cederbaum et al. in 1997 and termed intermolecular Coulombic decay (ICD).¹⁸ In ICD, the inner-valence vacancy is filled by an electron from an outer shell and

the energy released is transferred to a neighboring molecule, which is in turn ionized through the emission of a low-energy electron. This nonlocal decay mechanism is ultrafast and evolves usually on a femtosecond time scale.^{19–21} It results in the formation of two positive ions, leading typically to a Coulomb explosion of the system.

Experiments on ICD occurring after a manifold of excitation schemes have been identified in numerous systems,^{22,23} such as van der Waals clusters,^{24–28} hydrogen-bonded systems,^{29–33} endohedral fullerenes,³⁴ helium droplets,³⁵ benzene dimers,³⁶ liquid water,^{37,38} and condensed-phase interfaces.³⁹ The secondary ICD electrons usually have kinetic energies less than 10 eV and therefore have the capability to efficiently damage chemical bonds of neighboring molecules. Due to this key feature, ICD is expected to be a serious cause of DNA damage. Recent investigations on the hydrated biomolecules, where water is linked to biochemically relevant molecules via a hydrogen bond, imply that ICD can cause damage of biosystems directly by ionization.^{32,33} Therefore, studying this mechanism in multiple hydrogen-bond systems would be of fundamental importance for understanding DNA damage at the molecular level.

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Figure 1. Illustration of ICD process in a double hydrogen-bonded formic acid dimer. (a) The ICD process is initiated by an innervalence (iv) ionization of one of the FA molecules in the dimer. (b) The created iv vacancy is filled by an electron from the outer-valence (ov) shell at this site, and the energy released is transferred to the neighboring site from where a secondary low-energy electron is emitted. (c) Final fragmentation processes of the FA⁺...FA⁺ dicationic dimer with a Coulomb explosion pathway.

In the present work, we report experimental evidence of ICD in FAD initiated by electron impact ($E_0 = 90 \text{ eV}$). Here, the electron-initiated processes were thought to be responsible for an essential part of the radiation effects in gases and condensed matter.^{40–42} Below we study this decay process through fragment ion and electron coincident momentum spectroscopy, accompanied by high-level electronic structure calculations³² and *ab initio* molecular dynamics (AIMD) simulations. This would allow us to unambiguously identify the ionization mechanisms and elucidate the fragmentation dynamics of the FA⁺...FA⁺ dicationic dimer. Experiments were carried out using a multiparticle imaging spectrometer (reaction microscope) combined with a pulsed photoemission electron source.^{43–45} Our AIMD simulations of the FA⁺...FA⁺ dication were performed based on the Born–Oppenheimer approximation in which six different isomers of neutral FAD are considered⁴⁶ (see Methods).

We have computed the double ionization threshold of FAD, which amounts to 28.45 eV at the EOM-DIP-CCSD/cc-pVTZ level (see Methods). This corresponds to single ionization of both molecules within the dimer and is thus attributed to the ICD final state. Previous studies of the inner-valence ionization spectra of FA molecule^{47–50} have shown that the carbon 2s (C 2s) type states contribute mainly to the ionization-energy region of about 20–28 eV, while the ionization bands above 28 eV are mostly built on oxygen 2s (O2s) type states, i.e., the 5a' and 4a' orbitals. Hence, these states are energetically accessible for initiating the ICD process in FAD. More details will be provided in the following section. Here, the ionization and decay process is started by removal of an electron from the

O2s inner-valence shell of FA molecule upon electron impact (Figure 1a). Afterward the $FA^{+*}(O2s^{-1})\cdots FA$ dimer cation undergoes ICD, i.e., an electron from higher-lying orbital of $FA^{+*}(O2s^{-1})$ fills the O2s vacancy, and an electron from the neighboring FA molecule is ejected, forming the $FA^{+}\cdots FA^{+}$ dicationic dimer (Figure 1b). The repulsion between the two positive charges leads to Coulomb explosion (CE) of the system (Figure 1c). The present study aims to understand the mechanistic details of double hydrogen-bond breaking in FAD following ICD.

In our experiments, two fragment ions and one of the three outgoing electrons are detected in triple coincidence such that during offline analysis, the mass-over-charge ratios, the momentum vectors, and the kinetic energies for three charged particles are determined (see Methods). We identify the dimer breakup channels using a time-correlation map which is shown in Figure 2. After the ionization and decay process the dimer



Figure 2. Correlation of the time-of-flight (TOF) of the two measured cations. The distributions at the solid, dashed, and dotted diagonal lines show the breakup of the dicationic dimer into $FA^+ + FA^+$, $FA^+ + [FA - H]^+ + H$, and $[FA - H]^+ + [FA - H]^+ + 2H$ channels, respectively. The vertical and the horizontal dotted, dashed and solid lines represent the TOF of $[FA - H]^+$, FA^+ , and $[FA + H]^+$ ions, respectively. The counts intensity is color-coded on a linear scale with a color bar shown on the right.

fragments into two singly charged ions which are emitted backto-back with momenta of equal magnitude but opposite direction. As the time-of-flight (TOF) of an ion depends on its momentum and mass, the back-to-back emission leads to a diagonal line structure in the time-correlation map of both particles. The solid diagonal line indicates the position of events due to direct two-body fragmentation (FA⁺ + FA⁺). The dashed and dotted diagonal lines indicate the breakup channels where FA⁺ emitted a neutral hydrogen atom, i.e., FA⁺ + [FA – H]⁺ + H (dashed line), and double hydrogen-loss [FA – H]⁺ + [FA – H]⁺ + 2H (dotted line). The correlation lines become broader for these hydrogen-loss channels due to the missing momenta of the undetected neutral species.

In the following, we discuss only the $FA^+ + FA^+$ CE channel that appears as the solid diagonal line in Figure 2 and can clearly be assigned to ionization and two-body fragmentation recorded in coincidence with the FA⁺ + FA⁺ ion pair to show that ICD actually occurs for this fragmentation channel. The measurements of the scattered electrons allow us to determine the initial state of the decay process using the projectile energy loss (E_{loss}), as shown in Figure 3a,b, which is obtained by the



Figure 3. Measured and calculated energy spectra. (a) Projectile energy loss spectra for the inner-valence ionization of formic acid monomer that forms CH^+ and C^+ ions and for He(1s) ionization (solid line). (b) Energy loss spectrum for the double ionization of formic acid dimer leading to $FA^+ + FA^+$ ion pair. (c) Kinetic energy distributions of the emitted electrons for inner-valence ionization of FA monomer (solid line) and for ionization and fragmentation of FAD into $FA^+ + FA^+$. The spectra were normalized to each other in the energy range from 20 to 30 eV. (d) Difference spectrum obtained by subtracting the electron energy distribution of the monomer from the dimer shown in (c) and also ADC calculation³² (solid line). (e) Calculated single-ionization spectrum of the FAD. Note that a different intensity scale is used between 25 and 40 eV.

incident electron energy (E_0) minus the energy of the scattered projectile (E_1) . The minimum energy leading to specific ionization channels can be determined from the onsets of the E_{loss} spectrum, e.g., the E_{loss} of about 24.5 eV is obtained for the ionization of helium (He⁺); see the solid line in Figure 3a, which is measured for the energy calibration of the experimental spectra. Also included in Figure 3a is the E_{loss} spectrum corresponding to the inner-valence ionization of the FA monomer, which can result in the CH⁺ and C⁺ products with an E_{loss} of about 28.5 eV. The measured E_{loss} spectrum for the FA⁺ + FA⁺ fragmentation channel is presented in Figure 3b. This spectrum shows an onset at $E_{\rm loss} \sim 28.0-28.5$ eV in agreement with our EOM-DIP-CCSD calculations, which lie in the range of the inner-valence ionization energies of 5a' and 4a' molecular orbitals.⁴⁷⁻⁵⁰ Figure 3e presents the singleionization spectrum of FAD computed at ADC(3)/cc-pVTZ level (see Methods). It reveals that the removal of an O2s electron leads to the population of a large number of states that lie above the double-ionization threshold of FAD (\sim 28.0 eV). These states, spread over nearly 10 eV (see the zoomed-in part of the spectrum above 25 eV in Figure 3e), can all further decay by emitting an electron from the neighboring FA molecule in an ICD process.

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The ICD electrons are identified in Figure 3c,d which show the electron energy spectra up to 30 eV measured in coincidence with the $FA^+ + FA^+$ ion pair. The ICD electron emitted in Figure 1b cannot be seen separately since the spectrum contains background contribution of slow electrons arising from the initial direct electron-impact ionization process; see Figure 1a. In order to reveal the pure ICD electron spectra, we consider the electron energy spectrum for the inner-valence ionization of FA molecule as a reference (solid line) where ICD is absent, i.e., the spectrum corresponding to the $CH^{\scriptscriptstyle +}$ and $C^{\scriptscriptstyle +}$ products. This spectrum is normalized individually to the result of the dimer in the energy range above 20 eV where no ICD electrons contribute; see Figure 3c. An enhanced production of low-energy electrons is clearly observed in the spectrum of the dimer compared with the result of the monomer.^{27,45} The difference, as shown in Figure 3d, is the pure ICD electron energy spectrum for the $FA^+ + FA^+$ CE channel. This spectrum shows a kinetic energy peak at about 2.5 eV with an energy range up to about 7 eV. The energy range of this spectrum is consistent with our calculations, which show a broad band for the binding energies of the O2s⁻¹ states mainly located at 30-40 eV (Figure 3e and Figure 1 of the Supporting Information). The lowest doubly ionized state of the FA molecule lies at about 33.24 eV, which is computed at the EOM-DIP-CCSD/cc-pVTZ level. This means that the states whose binding energies are higher than 33 eV may also relax through intramolecular Auger decay and form a FA⁺⁺…FA.³² These results indicate that the singly ionized states lying at 28-33 eV are energetically accessible for the ICD channel in FAD, and thus, the ICD electron kinetic energies up to roughly 5 eV are expected for the $FA^+ + FA^+ CE$ channel. Furthermore, we performed calculations for the ICD electron spectrum (solid line in Figure 3d) using the approach described in ref 32 (see Figure 3 of the Supporting Information), which is in good agreement with the experimental result. We note that a low-energy peak (<2.0 eV) of the ICD electron spectrum is not seen in the measurements as the efficiency of the electron spectrometer quickly drops at this energy range.^{33,43}

In addition, as the obtained \vec{E}_{loss} for FA⁺ + FA⁺ ion pair is close to the double ionization threshold of FAD, one may expect that another fragmentation pathway, namely, sequential ionization (where the projectile electron successively kicks out one electron from each molecule of the dimer and thus leads to the repulsive FA⁺ + FA⁺ ion pair), can contribute to this CE channel. Our analysis on the absolute cross section, however, indicates that the sequential ionization is only a minor channel³³ due to its low cross section (see Methods).

The measured kinetic energy release (KER) spectrum for the FA⁺ + FA⁺ CE channel is presented in Figure 4a which shows a single-peak structure centered at KER ~ 4.15 eV. As the intermolecular potential drops with 1/R assuming two point-like charges separated by the distance *R*, we may deduce the Coulomb energy by considering that the two charges are located at the center-of-mass (COM) of the molecules. For a planar eight-membered ring structure of FAD (as shown in the inset in Figure 4a), which is suggested to be the lowest-energy configuration,⁴⁶ we determine the intermolecular COM



Figure 4. Kinetic energy release (KER) spectra of $FA^+ + FA^+$ ion pair. (a) Measured and *ab initio* molecular dynamics (AIMD) simulated KER spectra. Here, a planar eight-membered ring structure of FAD, corresponding to the global energy minimum geometry, was considered in the simulations. (b–f) AIMD simulated KER spectra corresponding to the isomers II–VI of FAD, respectively. The insets show the molecular geometries of different isomers of neutral FAD.⁴⁶

distance ($R_{\rm COM}$) of about 3.03 Å. This corresponds to KER ~ 4.75 eV for the complete fragmentation of the dimer into an FA⁺ + FA⁺ ion pair, which is about 0.6 eV higher than the measured KER value.

To quantify this CE process, we have performed AIMD simulations starting from the sampled neutral dimers with a given initial temperature (30 K) and instantaneous removal of the two outermost electrons from each of the molecules. This means that the simulations are started from the final state of ICD, i.e., FA⁺...FA⁺, corresponding to the electronic ground state of the dicationic dimer. We calculated the fragmentation dynamics of six different isomers of the neutral FAD⁴⁶ using the atom-centered density matrix propagation (ADMP) method with the long-ranged density-functional theory at the ω B97XD/cc-pVTZ level, which takes into account the internal and the rotational dynamics of the fragments (see Methods).

In the calculations, the KER value for each single trajectory is obtained by adding the Coulomb energy corresponding to the $R_{\rm COM}$ at t = 500 fs, which amounts to $R_{\rm COM} \sim 43$ au, to the KER resulting from the trajectory propagation.^{51,52} As shown in Figure 4a, the calculated KER considering a planar eightmembered ring structure of neutral FAD reproduces very well the experimental spectrum. We note that the present simulations assume an instantaneous ICD decay, neglecting the molecular dynamics of the singly ionized FA⁺*(O2s⁻¹)... FA dimer before the onset of ICD and the decay to the doubly ionized FA⁺...FA⁺ state. This may cause the slight deviations between calculated and measured KER in peak position (~0.1 eV) and the narrow peak width of the calculated KER spectrum. The calculated KER distributions for II-VI isomers are presented in Figures 4b-f, respectively, which show singlepeak structures centered at around 3.2-3.6 eV. The molecular structures for different isomers are also presented in the insets of Figure 4. These values are at least 0.5 eV smaller than the experimental and calculated KER shown in Figure 4a. This confirms the predominance of a planar eight-membered ring structure of the neutral FAD⁴⁶ in the present supersonic gas experiment. To estimate the part of energy that is transferred to internal and rotational motion of the fragments during the Coulomb explosion, we calculated the initial Coulomb energy (iCE) at t = 0, which corresponds to the Coulomb repulsion energy at R_{COM}, and assumed that all this energy is then entirely released as translational kinetic energy of the fragments. The results are shown in Figure 2 of the Supporting Information. The iCE spectrum corresponding to the planner eight-member ring isomer shows a peak at higher KER (~4.75 eV), indicating that the amount of Coulomb energy deposited into internal and rotational energy of the two FA⁺ fragments is roughly 0.7 eV.

To uncover the underlying dynamics of the present CE channel, we obtain the detailed temporal evolutions of the $FA^+\cdots FA^+$ dications with the predominant structure of FAD, which are presented in Figure 5. These results allow us to



Figure 5. Fragmentation dynamics of the dicationic dimer into $FA^+ + FA^+$ pathway. (a) AIMD simulated center-of-mass (COM) distance (R_{COM}) between two FA⁺ cations as a function of propagation time. (b, c) Path-integral trajectories in the time range of 0–25 fs (b) and 25–50 fs (c). (d–f) AIMD simulated KER of FA⁺ + FA⁺ ion pair (d) and vibrational energy E_{vib} (e) and rotational energy E_{rot} (f) of FA⁺ cations as a function of time. The mean values for E_{vib} and E_{rot} are shown by the solid lines in (e) and (f), respectively. The color bar is linear with arbitrary units.

elucidate the mechanistic details responsible for the double hydrogen-bond breaking and how the energy is distributed among the translational, vibrational, and rotational degrees of freedom of the two FA⁺ molecular cations during the explosion. It can be seen from Figure 5a and Figure 5d that after the decay process the two FA⁺ ions fly apart rapidly under the Coulomb repulsive force, i.e., both R_{COM} and KER of the FA⁺ + FA⁺ ion pair are starting to increase once ICD occurs.

The vibrational and rotational motions of fragments are investigated by deducing the vibrational and rotational energies for each of the FA⁺ cations from the AIMD simulations.⁵ The vibrational energy (E_{vib}) of FA⁺ is presented in Figure 5e as a function of time, which shows a constant energy-spread evolution centered at $E_{\rm vib} \sim 0.53$ eV over the whole range of propagation time (0-500 fs), indicating that the partitioning of KER and vibrational energy in the molecular fragments is determined early in the CE such that the total vibrational energy is constant over these time scales. While for the rotational energy (E_{rot}) of FA⁺, shown in Figure 5f, we observe a fast increase of E_{rot} , due to the intermolecular Coulomb force, that converges to a constant value ($\sim 125 \text{ meV}$) at larger time (>100 fs). This increased rotational frequency would cause the hydrogen-bond breaking of the system. It can be seen from the path-integral trajectories (0-50 fs) of the FA⁺ + FA⁺ ion pair, shown in Figure 5b,c, that the two hydrogens involved in the hydrogen bonding rotate very quickly out of the coplanar ring of the dimer. This leads to a fast and concerted breaking of the double hydrogen-bond in the system, while the molecular separation between two FA⁺ is only slightly increased by the Coulomb force during this time.

Ultrafast intermolecular decay processes in a double hydrogen-bond dimer of formic acid (FA) molecules have been investigated using fragment ions and electron coincident momentum spectroscopy and also ab initio calculations. We observe direct evidence of ICD in the decay of formic acid dimer (FAD) ionized by electron impact ($E_0 = 90$ eV). First, the back-to-back emission of ion pairs reveals a two-body dissociation property of the fragmenting dicationic dimers. Second, the measurements of projectile energy loss distributions in coincidence with the ion pair indicate that the reaction is initiated by removal of an electron from the inner-valence (oxygen 2s) shell of the FA molecule, while a sequential process in which the projectile ionizes one molecule after the other is only a minor channel here. Third, the measured electron spectra provide an important signature of the decay process, i.e., the enhancement of low-energy (<10 eV) ICD electrons in comparison with the ionization spectrum of isolated FA molecules. Furthermore, the intermolecular proton transfer is not identified in the present experiments, ^{12,16,53} e.g., $[FA - H]^+ + [FA + H]^+$ channel, supporting that ICD is a very efficient process taking place before the proton transfer.

This ultrafast mechanism forms a dicationic dimer with two charges located at each of the two molecules, leading typically to a Coulomb explosion of the system. The fragmentation dynamics of $FA^+ \cdots FA^+$ dications were interpreted with the help of AIMD simulations showing that the measured KER spectrum of $FA^+ + FA^+$ ion pair is very well reproduced only when starting from a planar eight-membered ring structure of the neutral dimer, while the other isomers are found to be minor contributions to the present supersonic gas experiment. Furthermore, our study reveals a fast and concerted breaking of the double hydrogen-bond initiated by the strong rotational motions of two FA^+ cations following the intermolecular decay mechanism and thus discloses a potential role of ultrafast molecular rotation in breaking the intermolecular hydrogenbonds. The present study of ICD in formic acid dimers is expected to be relevant in radiation damage of biological tissues, both for the fundamental reason for generating genotoxic particles (low-energy electrons and radicals) and because of the double hydrogen-bond breakage induced by this ultrafast decay process. Indeed, our further calculations indicate that it may also occur in the hydrogen-bonding adenine—thymine and guanine—cytosine base pairs, where the computed double ionization thresholds are 18.6 and 18.3 eV, respectively, at the CCSD(T)/cc-pVDZ level. These threshold values are much lower than for FAD, suggesting that ICD is energetically accessible for DNA base pairs.

METHODS

Experimental Methods. The experimental data were obtained by crossing an electron projectile beam with a gas target jet and employing a multiparticle coincidence spectrometer (reaction microscope) for detection of the charged collision fragments.^{43,45} The formic acid dimers are generated in a supersonic expansion of helium gas (stagnation pressure 1 bar) with seeded formic acid vapor. The pulsed electron beam is produced in an electron gun from a tantalum photocathode which is irradiated by UV-light pulses of 0.5 ns duration, and it is collimated by electrostatic lens elements. The energy width of the electron beam is about 0.5 eV.⁴⁴ It is guided by an axial magnetic field (0.7 mT) to the crossing zone with the gas jet and further to a beam dump which is a central bore in the electron detector. Secondary electrons and ions are extracted by means of a homogeneous electric field to opposite directions and projected onto two position- and time-sensitive multichannel plate detectors with 80 mm diameter of the active area. From the impact positions and the times-of-flight the momentum vectors and consequently the kinetic energies of the particles emerging from the reaction are determined. The acceptance angle for detection of low energy electrons up to the kinetic energy of 15 eV is almost 4π with the exception of small forward and backward angles which are lost due to the detector bore. In order to maximize the acceptance for molecular ion fragments, the electric extraction field of 1.0 V cm⁻¹ is ramped up to 20 V cm⁻¹ after 400 ns when the electrons have reached the detector. In our experiment, monomer ionization is simultaneously recorded with cluster ionization.

The absolute cross section for SI leading to $FA^+ + FA^+$ ion pair is calculated by multiplying the partial ionization cross section for the first collision partner with the probability to subsequently ionize the second molecule. The latter is the relative size of its ionization cross section compared to the total surface area of a sphere with radius of the intermolecular distance. The calculations can be expressed as following:

$$\sigma_{\rm SI}(E_0) = \sigma_{\rm FA}^+(E_0) \frac{\sigma_{\rm FA}^+(E_1)}{4\pi R^2}$$
(1)

Here $\sigma_{FA}^+(E_0 = 90 \text{ eV}) \approx \sigma_{FA}^+(E_1 = 78.5 \text{ eV}) \approx 1.0 \text{ Å}^2$ is the partial ionization cross section for production of FA⁺ at the projectile energy of $E_0 = 90 \text{ eV}$ and $E_1 = 78.5 \text{ eV}$ which considers the minimum energy loss in the first collision.⁵⁴ R = 3.0 Å is the intermolecular center-of-mass distance obtained from the molecular geometry calculations.⁴⁶ As a result, we obtain $\sigma_{SI}(90 \text{ eV}) \approx 0.885 \times 10^{-18} \text{ cm}^2$. The oxygen 2s (O2s)

ionization cross section of FA molecule is not available in the literature, we refer to a value of about 1.73×10^{-17} cm² for the O2s ionization of water.³³ This indicates that SI in the formic acid dimer is only a minor channel in comparison to O2s inner-valence ionization of molecules.

Computational Methods. The double ionization thresholds of formic acid dimer (28.45 eV) and monomer (33.24 eV) were obtained at the EOM-DIP-CCSD/cc-pVDZ level of theory with the nonelectron-conserving variant of EOM-CCSD where target states are reached by detaching two electrons from the corresponding ground state.^{55,56} The calculations were performed using the Q-Chem 4.2 program suite.⁵⁷

We performed calculations for the fragmentation dynamics of the FA⁺…FA⁺ dication on the electronic ground state. In these calculations, we considered six lowest energy isomers of the formic acid (FA) dimers (Figure 4) including the planar eight-membered ring structure that corresponds to the global energy minimum geometry.⁴⁶ The AIMD simulations were performed in two steps: (i) the initial conditions, i.e., geometries and velocities of every atom, of neutral (FA)₂ were sampled by the quasi-classical fixed normal-mode sampling method⁵⁸ under the temperature of about 30 K in which the populations of the initial vibrational states were determined by Boltzmann distributions; (ii) we assumed an instantaneous vertical transition to the electronic ground state of the doubly charged FA⁺ ··· FA⁺ dimer. The molecular dynamics simulation was performed under the extended Lagrangian molecular dynamics scheme adopting the so-called atom-centered density matrix propagation (ADMP) method⁵⁸⁻⁶⁰ using the long-ranged density-functional theory method at the ω B97XD/cc-pVTZ level. We obtained the dynamical simulations for 500 fs time range with 0.5 fs time step under the approximate self-consistent field convergence at each time step. This method enables Born-Oppenheimer molecular dynamics on a given electronic state with accurate long-term energy conservation. All calculations were carried out with the Gaussian 16 suite of programs.⁶¹

The vertical ionization spectrum of FAD was computed using the third-order algebraic diagrammatic construction (ADC(3)) approximation scheme for the one-particle Green's function. 62-65 The ADC(3) method describes the energies and pole strengths of ionization transitions involving the "one-hole" (1h) and "two-hole–one-particle" (2h-1p) electronic states through third and first order, respectively, in the residual electronic interaction. The ADC(3) eigenstate equations⁶⁶ are solved using the Block–Lanczos procedure which is applicable in situations where the breakdown of the orbital picture of ionization⁶⁷ takes place. The latter phenomenon manifests itself by a strong redistribution of spectral intensity from the 1h main states to 2h-1p satellites and is typical for inner-valence ionization considered in this work. The ADC(3) method was previously applied in various molecules including computations of ionization spectra of the valence,^{68–70} inner-valence,⁷¹ and K-shell⁷² orbitals and is proved to be an accurate and valuable tool in spectroscopic studies. The cc-pVTZ basis set was used in the calculations. The resulting spectral envelopes (constructed by convolving the generated spectra with Gaussians of the fwhm (full width at half-maximum) equal to 0.2 eV) are shown in Figure 3e together with the underlying transitions, which are presented by the black lines. The contributions of the inner-valence O2s holes are presented by the colored envelopes in the energy region from 25 to 40 eV.

ASSOCIATED CONTENT

Supporting Information

The Supporting Information is available free of charge at https://pubs.acs.org/doi/10.1021/acs.jpclett.2c00957.

Detailed description of the single ionization spectra of FAD in the $O2s^{-1}$ energy region, calculated initial Coulomb energy of FA⁺ + FA⁺ channel, and calculated ICD electron spectrum (PDF)

Animated movie for the Coulomb explosion process (AVI)

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Notes

The authors declare no competing financial interest.

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